Alfven waves: small-scale turbulence and large-scale structure

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MHD turbulence: Alfvenic cascade $\partial \mathbf{z}^{\pm} \mp (\mathbf{v}_A \cdot \nabla) \mathbf{z}^{\pm} + (\mathbf{z}^{\mp} \cdot \nabla) \mathbf{z}^{\pm} = -\nabla P + \frac{1}{R_e} \nabla^2 \mathbf{z}^{\pm} + \mathbf{f}^{\pm}$ Ideal system conserves the Elsasser energies $E = \frac{1}{2} \int (v^2 + b^2) d^3x$ $= E = \frac{1}{2} \int (v^2 + v^2) a \\ H^C = \int (\mathbf{v} \cdot \mathbf{b}) d^3 x$ $E^+ = \int (\mathbf{z}^+)^2 d^3 x$ $E^- = \int (\mathbf{z}^-)^2 \, d^3 x$ \mathbf{B}_0 V۸ \mathbf{V}_A Z+ \mathbf{B}_{0} V_A \mathbf{V}_A $E^+ \gg E^-$: imbalanced case $E^+ \sim E^-$: balanced case.

$$H^{C} = \int (\mathbf{v} \cdot \mathbf{b}) d^{3}x = \frac{1}{4} (E^{+} - E^{-}) \neq 0$$
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Strength of interaction in MHD turbulence

$$\partial \mathbf{z}^{\pm} \mp (\mathbf{v}_{A} \cdot \nabla) \mathbf{z}^{\pm} + (\mathbf{z}^{\mp} \cdot \nabla) \mathbf{z}^{\pm} = -\nabla P + \frac{1}{Re} \nabla^{2} \mathbf{z}^{\pm} + \mathbf{f}^{\pm}$$

$$\underbrace{\langle k_{\parallel} v_{A} \rangle z^{\pm}}_{(k_{\perp} z^{\mp}) z^{\pm}} \underbrace{\langle k_{\perp} z^{\mp} \rangle z^{\pm}}_{(k_{\perp} z^{\mp}) z^{\pm}}$$

When
$$~~k_\parallel v_A \gg k_\perp z^\mp~$$
 turbulence is weak

When $~~k_\parallel v_A \sim k_\perp z^\mp$ turbulence is strong

Wave MHD turbulence: Phenomenology

Three-wave interaction of shear-Alfven waves

$$\omega(k) = |k_z| v_A$$

$$\begin{cases} \omega(k) = \omega(k_1) + \omega(k_2) \\ \mathbf{k} = \mathbf{k}_1 + \mathbf{k}_2 \end{cases}$$

Only counter-propagating waves interact, therefore, k_{1z} and k_{2z} should have opposite signs.

Either $k_{1z}=0$ or $k_{2z}=0$

Wave interactions change k_{\perp} but not k_z

At large k_1:
$$E(k_z,k_\perp) \propto g(k_z) k_\perp^{-eta}$$

Montgomery & Turner 1981, Shebalin et al 1983

Analytic framework [Galtier, Nazarenko, Newell, Pouquet, 2000]

In the zeroth approximation, waves are not interacting. and z^+ and z^- are independent:

$$\langle \mathbf{z}^{+}(\mathbf{k}) \cdot \mathbf{z}^{+}(\mathbf{k}') \rangle = e^{+}(k_{z}, k_{\perp}) \delta(\mathbf{k} + \mathbf{k}')$$

$$\langle \mathbf{z}^{-}(\mathbf{k}) \cdot \mathbf{z}^{-}(\mathbf{k}') \rangle = e^{-}(k_{z}, k_{\perp}) \delta(\mathbf{k} + \mathbf{k}')$$

$$\langle \mathbf{z}^{+}(\mathbf{k}) \cdot \mathbf{z}^{-}(\mathbf{k}') \rangle = 0$$

When the interaction is switched on, the energies slowly change with time: $e^{\pm}(k_z, k_{\perp}, t)$

$$\partial_t \mathbf{z}^{\pm} - (\mathbf{v}_A \cdot \nabla) \mathbf{z}^{\pm} + (\mathbf{z}^{\mp} \cdot \nabla) \mathbf{z}^{\pm} = -\nabla P$$

$$\partial_t \langle z^+ z^+ \rangle = \dots \langle z^- z^+ z^+ \rangle + \langle z^+ z^- z^+ \rangle \dots$$

$$\partial_t \langle z^- z^+ z^+ \rangle = \dots \langle z^+ z^- z^+ z^+ \rangle + \langle z^- z^- z^+ z^+ \rangle + \langle z^- z^+ z^- z^+ \rangle \dots$$

split into pair-wise correlators using Gaussian rule

Weak turbulence: Analytic framework [Galtier, Nazarenko, Newell, Pouquet, 2000]

$$\partial_t \langle z^+ z^+ \rangle = \dots \langle z^- z^+ z^+ \rangle + \langle z^+ z^- z^+ \rangle \dots$$

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split into pair-wise correlators using Gaussian rule

$$\partial_t e^{\pm}(k_z, k_\perp) = \int M_{k,pq} e^{\mp}(0, q_\perp) \left[e^{\pm}(k_z, k_\perp) - e^{\pm}(k_z, p_\perp) \right] \delta(\mathbf{k}_\perp - \mathbf{p}_\perp - \mathbf{q}_\perp) d^2 p d^2 q$$
$$M_{k,pq} = \frac{\pi}{v_A} \frac{(\mathbf{k}_\perp \times \mathbf{q}_\perp)^2 (\mathbf{k}_\perp \cdot \mathbf{p}_\perp)^2}{k_\perp^2 p_\perp^2 q_\perp^2}$$

This kinetic equation has all the properties discussed in the phenomenology: it is scale invariant, z^+ interacts only with z^- , k_z does not change during interactions. Weak turbulence: Analytic framework [Galtier, Nazarenko, Newell, Pouquet, 2000]

$$\partial_t e^{\pm}(k_z, k_{\perp}) = \int M_{k,pq} e^{\mp}(0, q_{\perp}) \left[e^{\pm}(k_z, k_{\perp}) - e^{\pm}(k_z, p_{\perp}) \right] \delta(\mathbf{k}_{\perp} - \mathbf{p}_{\perp} - \mathbf{q}_{\perp}) d^2 p d^2 q$$

Statistically balanced case: $e^+ = e^-$

$$e^+(k_z, k_\perp) = e^-(k_z, k_\perp) = g(k_z)k_\perp^{-3}$$

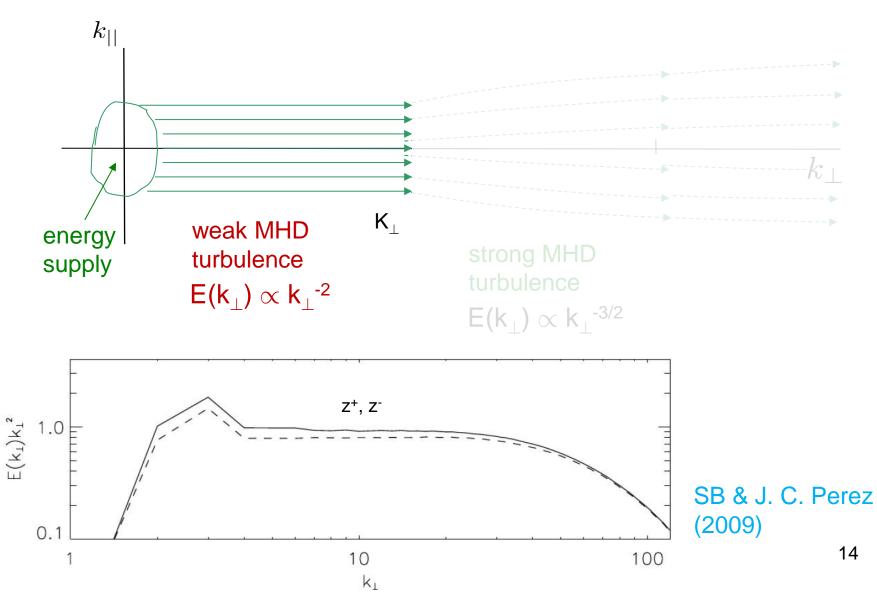
where $g(k_z)$ is an arbitrary function.

The spectrum of weak balanced MHD turbulence is therefore:

$$E^{\pm}(k_z,k_{\perp}) = e^{\pm}(k_z,k_{\perp})2\pi k_{\perp} \propto k_{\perp}^{-2}$$

Ng & Bhattacharjee 1996, Goldreich & Sridhar 1997

Weak MHD turbulence



Imbalanced weak MHD turbulence (where problems begin)

 $\partial_t e^{\pm}(k_z, k_{\perp}) = \int M_{k,pq} e^{\mp}(0, q_{\perp}) \left[e^{\pm}(k_z, k_{\perp}) - e^{\pm}(k_z, p_{\perp}) \right] \delta(\mathbf{k}_{\perp} - \mathbf{p}_{\perp} - \mathbf{q}_{\perp}) d^2 p d^2 q$

The kinetic equation has a one-parameter family of solutions:

$$\begin{aligned} e^+(k_z, k_\perp) &= g^+(k_z) k_\perp^{-3-\alpha} \\ e^-(k_z, k_\perp) &= g^-(k_z) k_\perp^{-3+\alpha} \end{aligned} \quad \text{with -1} < \alpha < 1 \end{aligned}$$

What do these solutions mean? Hint: calculate energy fluxes.

Assume that e+ has the steeper spectrum and denote the energy fluxes ϵ^+ and ϵ^- . Then $\epsilon^+ > \epsilon^-$

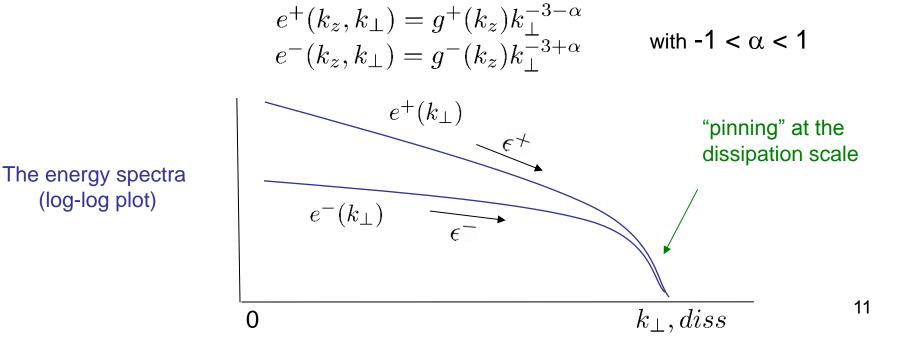
and:
$$\alpha = f(\epsilon^+/\epsilon^-)$$

Imbalanced weak MHD turbulence

(where problems begin)

$$\partial_t e^{\pm}(k_z, k_{\perp}) = \int M_{k,pq} e^{\mp}(0, q_{\perp}) \left[e^{\pm}(k_z, k_{\perp}) - e^{\pm}(k_z, p_{\perp}) \right] \delta(\mathbf{k}_{\perp} - \mathbf{p}_{\perp} - \mathbf{q}_{\perp}) d^2 p d^2 q$$

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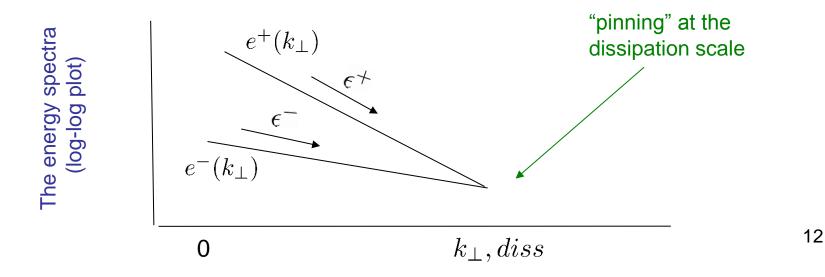
Imbalanced weak MHD turbulence (where problems begin)

$$e^{+}(k_{z}, k_{\perp}) = g^{+}(k_{z})k_{\perp}^{-3-\alpha} -1 < \alpha < 1$$

$$e^{-}(k_{z}, k_{\perp}) = g^{-}(k_{z})k_{\perp}^{-3+\alpha} \qquad \alpha = f(\epsilon^{+}/\epsilon^{-})$$

The spectra are "pinned" at the dissipation scale.

• If the ratio of the energy fluxes is specified, then the slopes are specified, but the amplitudes depend on the dissipation scale, or on the Re number.



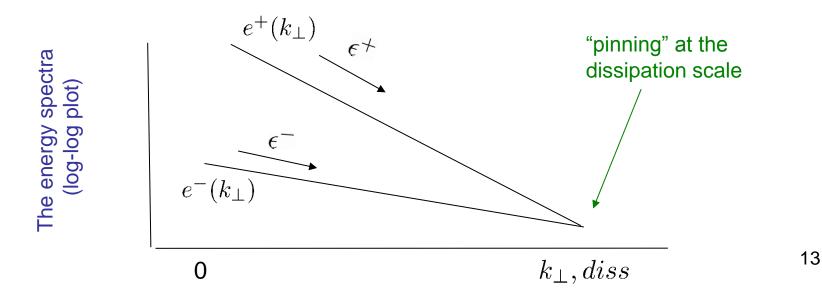
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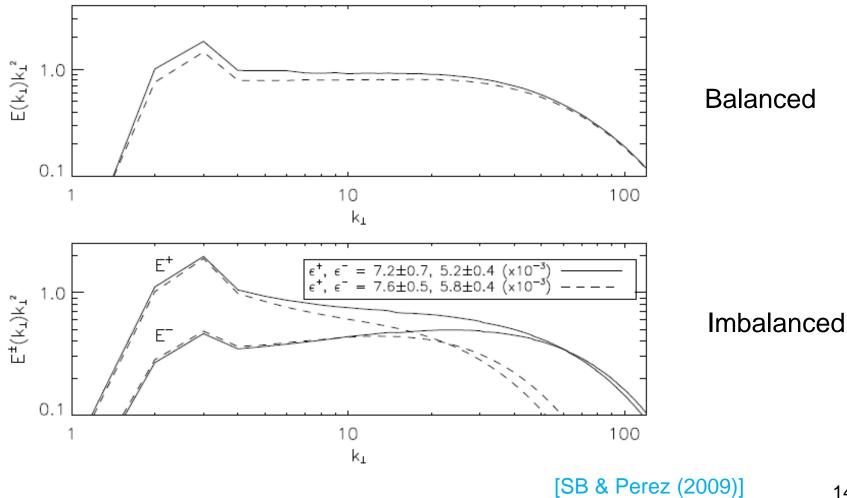
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Imbalanced weak MHD turbulence: Numerical results



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Residual energy in weak MHD turbulence

$$\begin{aligned} \langle \mathbf{z}^{+}(\mathbf{k}) \cdot \mathbf{z}^{+}(\mathbf{k}') \rangle &= e^{+}(k_{\parallel}, k_{\perp}) \delta(\mathbf{k} + \mathbf{k}') & \checkmark \\ \langle \mathbf{z}^{-}(\mathbf{k}) \cdot \mathbf{z}^{-}(\mathbf{k}') \rangle &= e^{-}(k_{\parallel}, k_{\perp}) \delta(\mathbf{k} + \mathbf{k}') & \checkmark \\ \langle \mathbf{z}^{+}(\mathbf{k}) \cdot \mathbf{z}^{-}(\mathbf{k}') \rangle &= q^{r}(k_{\parallel}, k_{\perp}) \delta(\mathbf{k} + \mathbf{k}') & \neq \mathbf{0} \\ \langle \mathbf{z}^{+} \cdot \mathbf{z}^{-} \rangle &= \langle v^{2} - b^{2} \rangle & \text{since the waves are not independent!} \end{aligned}$$

What is the equation for the residual energy?

SB & Perez PRL 2009

Residual energy in weak MHD turbulence

- Waves are almost independent one would not expect any residual energy!
- Analytically tractable:

$$\partial_t q^r = 2ik_{\parallel} v_A q^r - \gamma_k q^r +$$

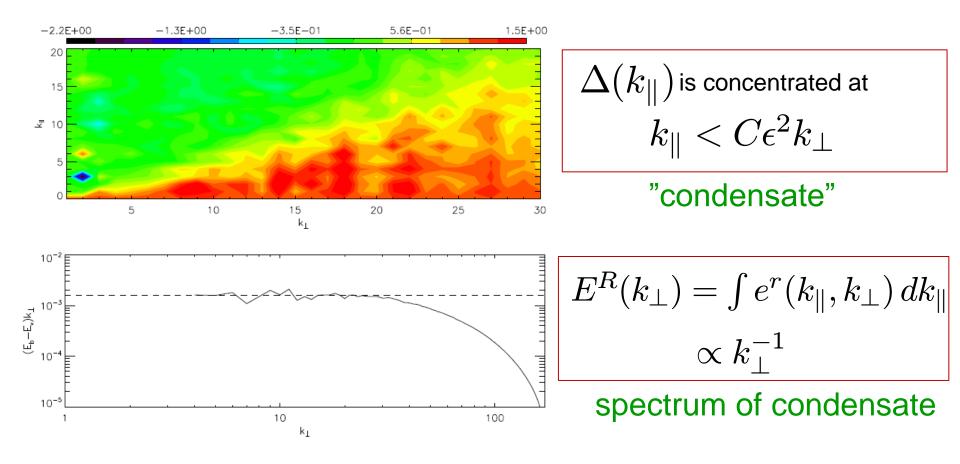
where:
$$R_{k,pq} = (\pi v_A/2)(\mathbf{k}_\perp \times \mathbf{q}_\perp)^2 (\mathbf{k}_\perp \cdot \mathbf{p}_\perp) (\mathbf{k}_\perp \cdot \mathbf{q}_\perp) / (k_\perp^2 p_\perp^2 q_\perp^2)$$

Conclusions:

- Residual energy is always generated by interacting waves!
- ∫ ... < 0, so the residual energy is negative: magnetic energy dominates!

Y. Wang, S. B. & J. C. Perez (2011) S.B, J. C. Perez & V. Zhdankin (2011)

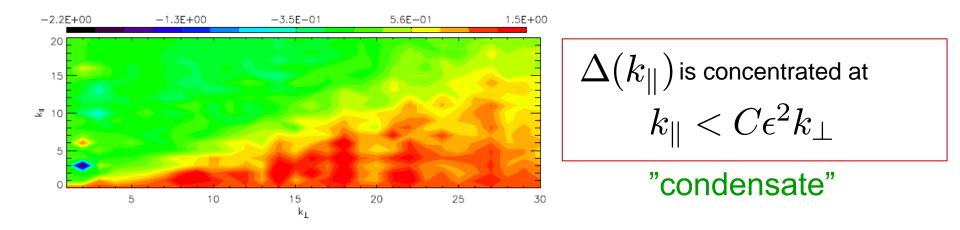
Residual energy in weak MHD turbulence $e^{r}(k) = Re\langle z^{+}(k) \cdot z^{-}(k) \rangle \propto -\epsilon^{2} k_{\perp}^{-2} \Delta(k_{\parallel})$

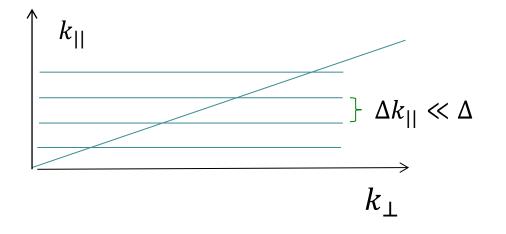


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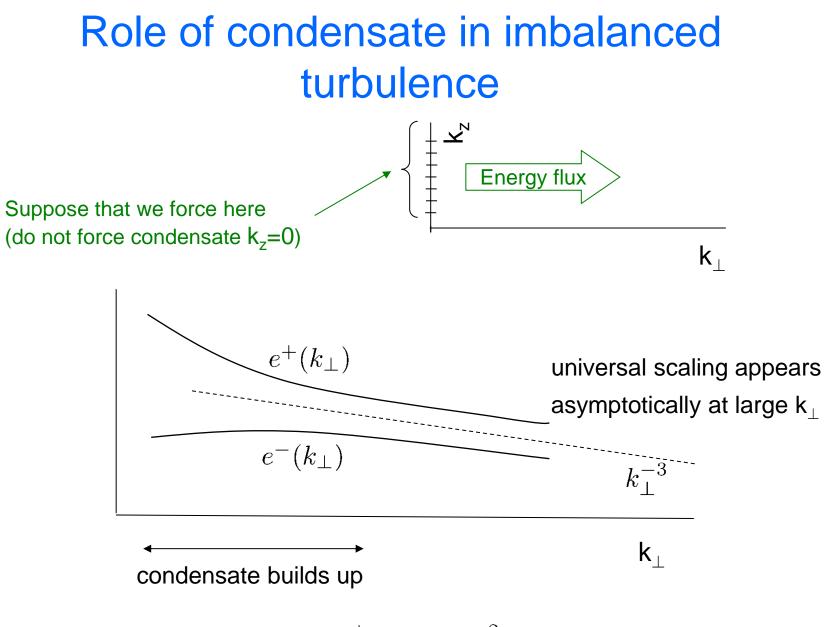
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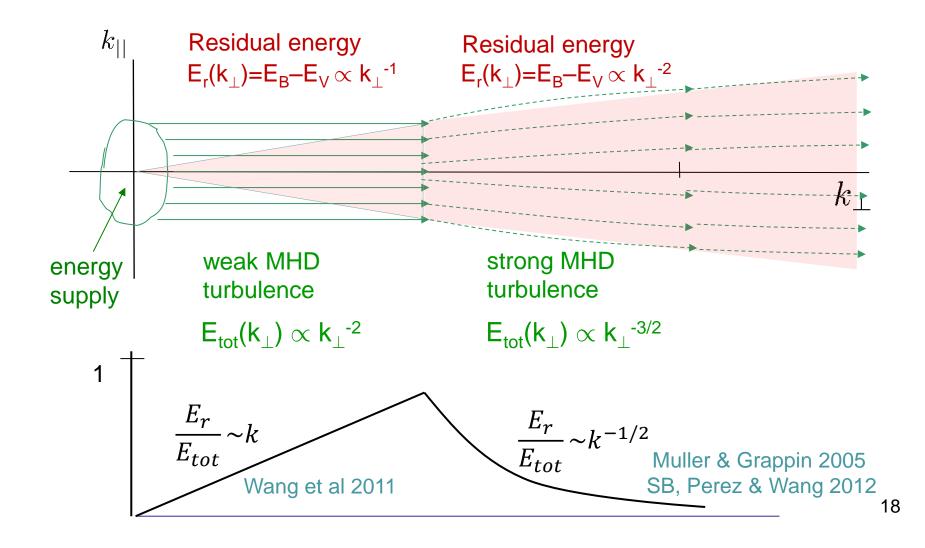


Structure in k-space must be resolved



$$E^{\pm}(k_{\perp}) \propto k_{\perp}^{-2}$$

Universal picture of MHD turbulence



Summary

- Weakly interacting Alfven waves spontaneously generates a condensate of energy (real and residual) at small $k_{||}$.
- The condensate a large-scale slowly evolving structure -- is a strongly nonlinear object. Waves are scattered by this structure. The wave spectrum is

$$E^{\pm}(k_{\perp}) \propto k_{\perp}^{-2}$$

• Especially important for imbalanced turbulence.