

# Self-Gravitating Supersonic MHD Turbulence

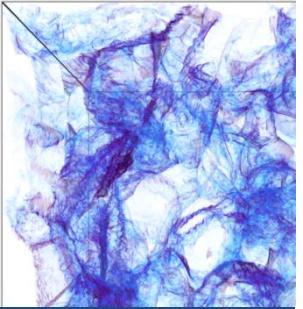
## GTP Workshop on MHD LES, Session VI



### Wolfram Schmidt

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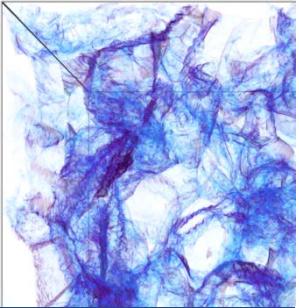
Wolfram Schmidt

Universität Göttingen

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# Application I: Star-Forming Clouds

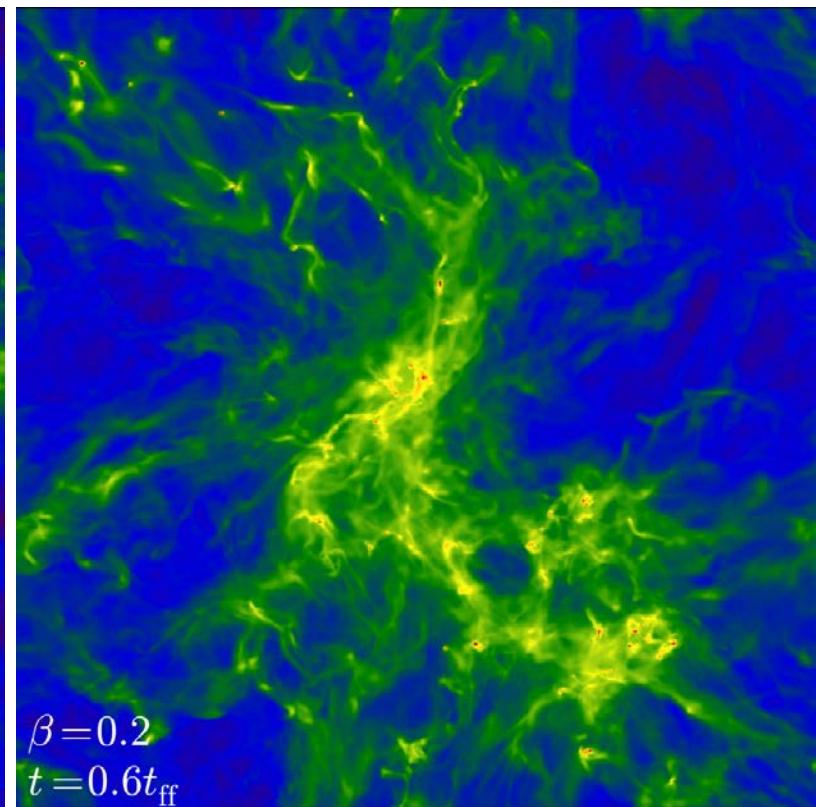
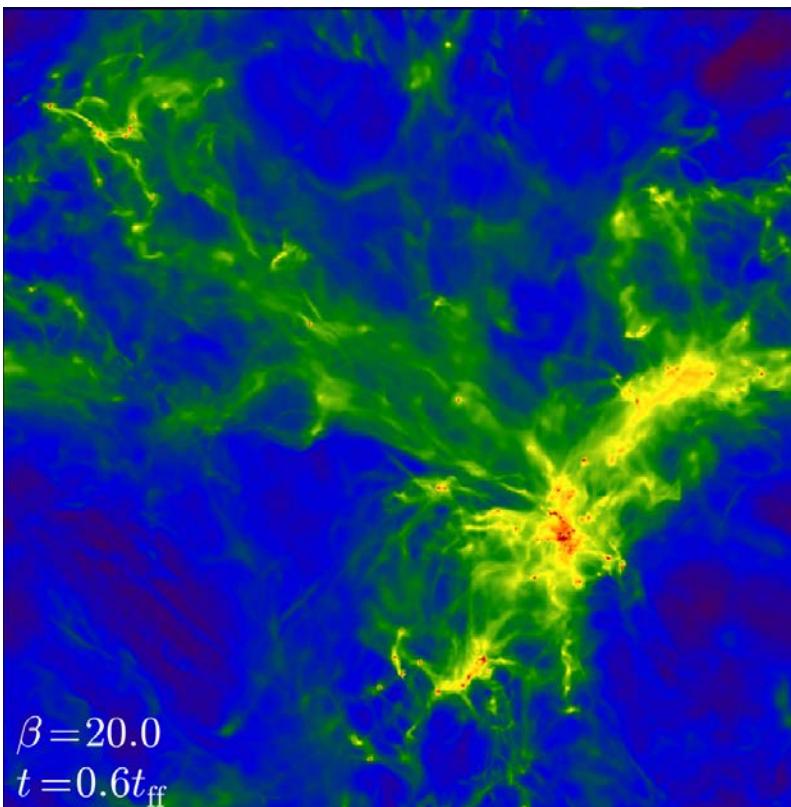


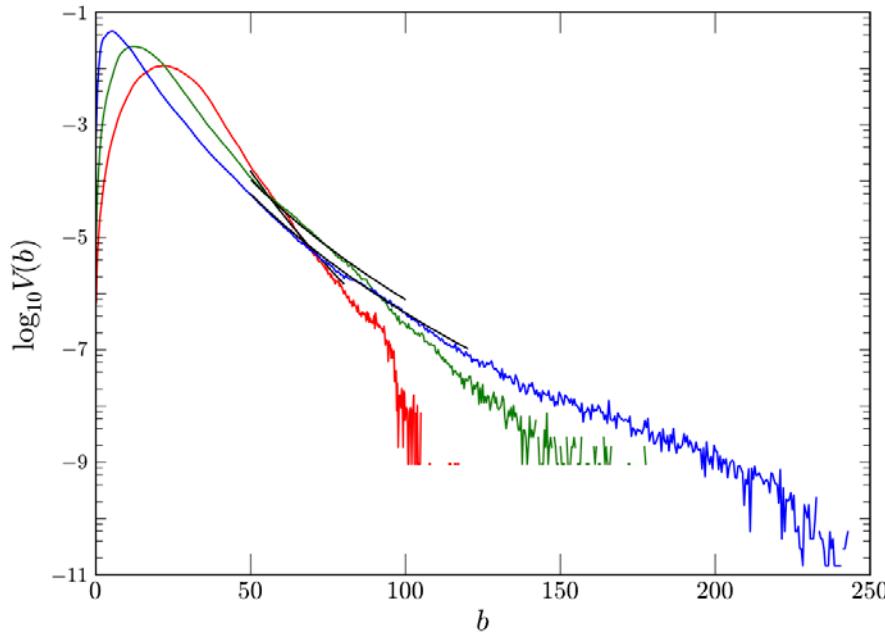
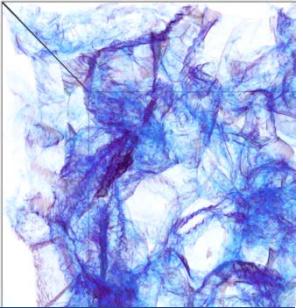


# ILES of Forced MHD

Collins et al. (2012):

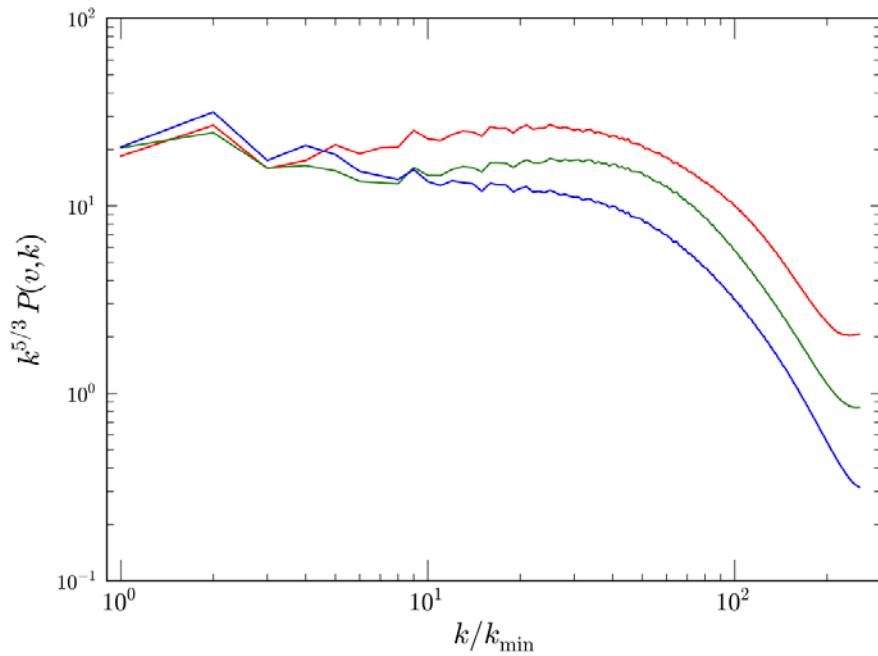
- AMR simulation (4 levels, 16 cells per Jeans length)
- PLM ([Li et al. 2008](#)), CT method,  $\text{Ma}_{\text{rms}} \approx 9$





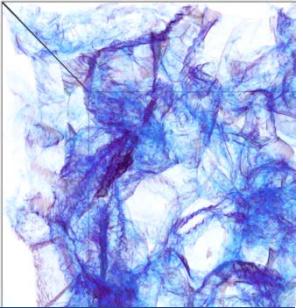
## Velocity power spectra:

- no evolution due to gravity
- Steeper slope for lower  $\beta_0$



**Magnetic field PDFs:**

- field fluctuation  $b$
- peaks given by  $\beta_0$
- widest tail for weakest field

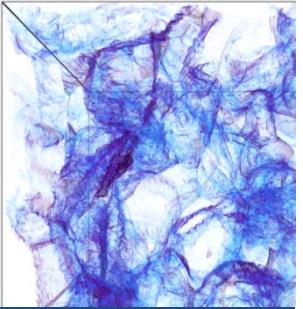


# Local Magnetic Field Amplification

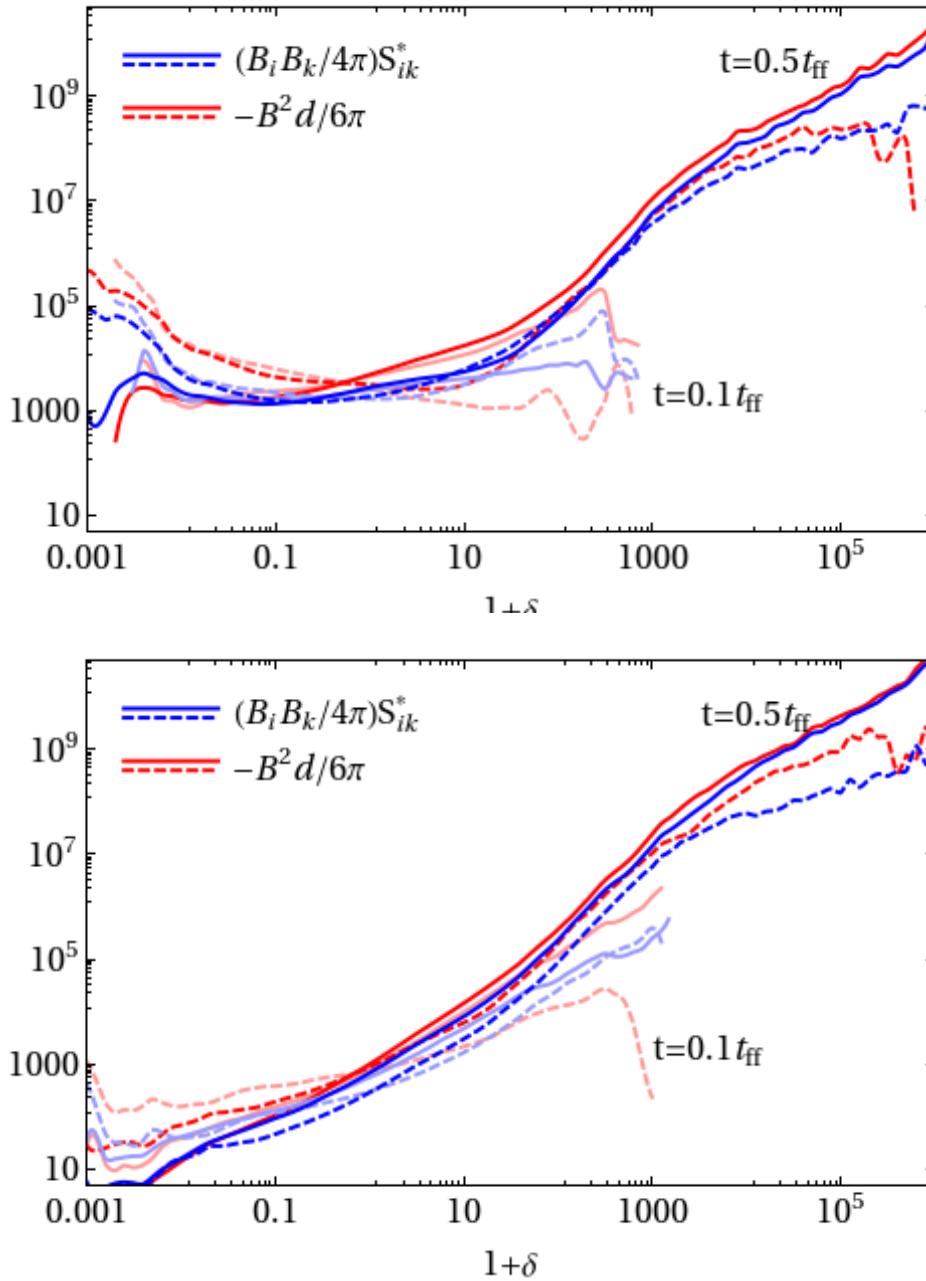
- PDE for the **magnetic pressure** follows from compressible induction equation (ideal MHD):

$$\frac{D}{Dt} \left( \frac{B^2}{8\pi} \right) = \frac{1}{4\pi} \left( B_i B_k S_{ik}^* - \frac{2}{3} B^2 d \right)$$

- Two contributions:
  - Amplification by **shear** (dynamo action)
  - Gravitational and shock **compression**
- Can be positive or negative
  - Compute averages of positive and negative contributions for given overdensity  $\rho / \rho_0 = 1 + \delta$

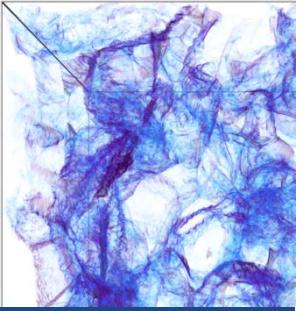


$\beta \approx 0.034$  and  $0.21$  (WS, Collins, Kritsuk 2013)



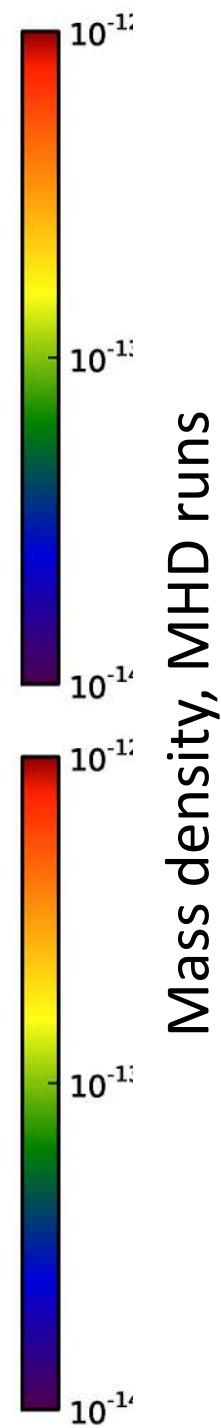
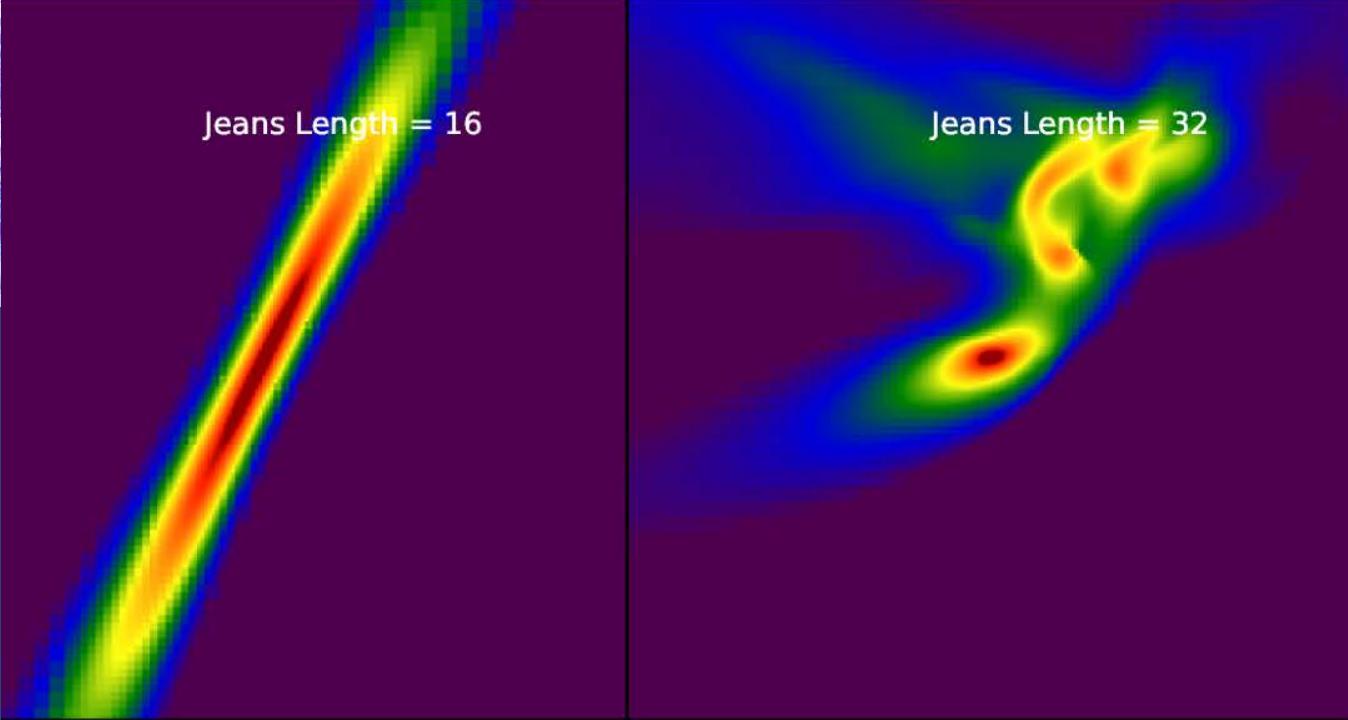
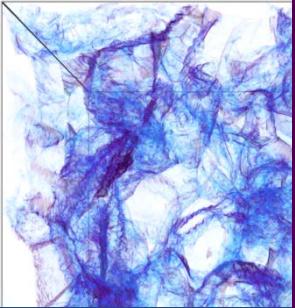
## Magnetic field amplification:

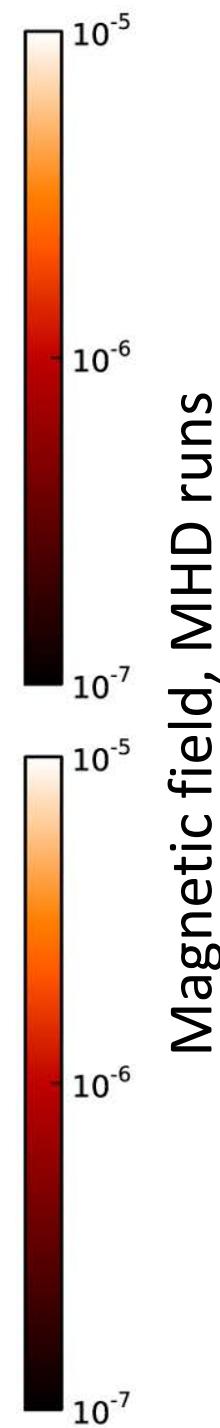
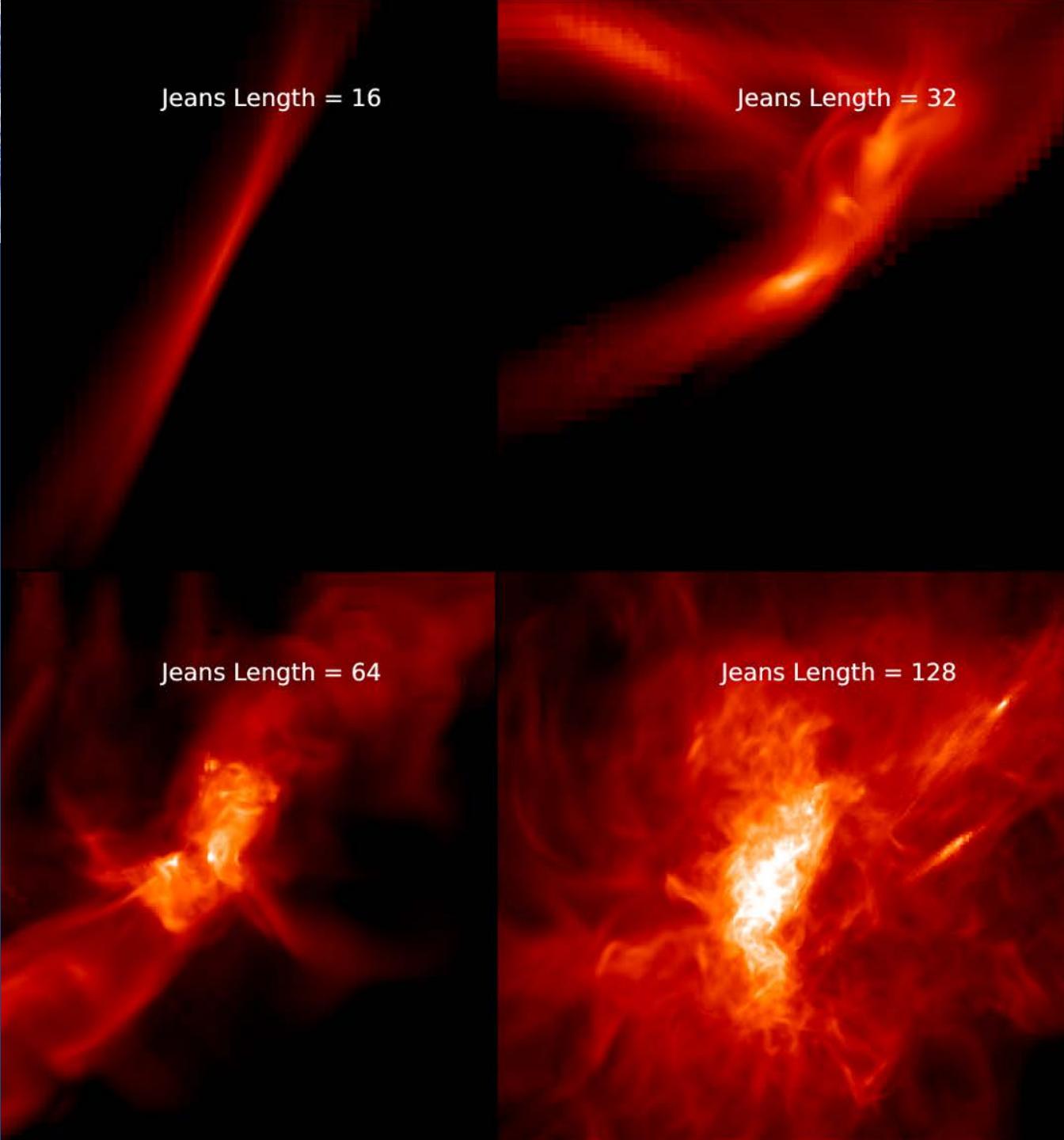
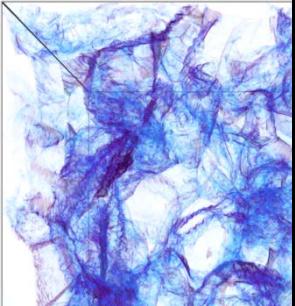
- initially saturation
- balance between shear-induced and compressive amplification
- stronger effect for higher  $\beta$
- net amplification at high densities (collapsing gas)

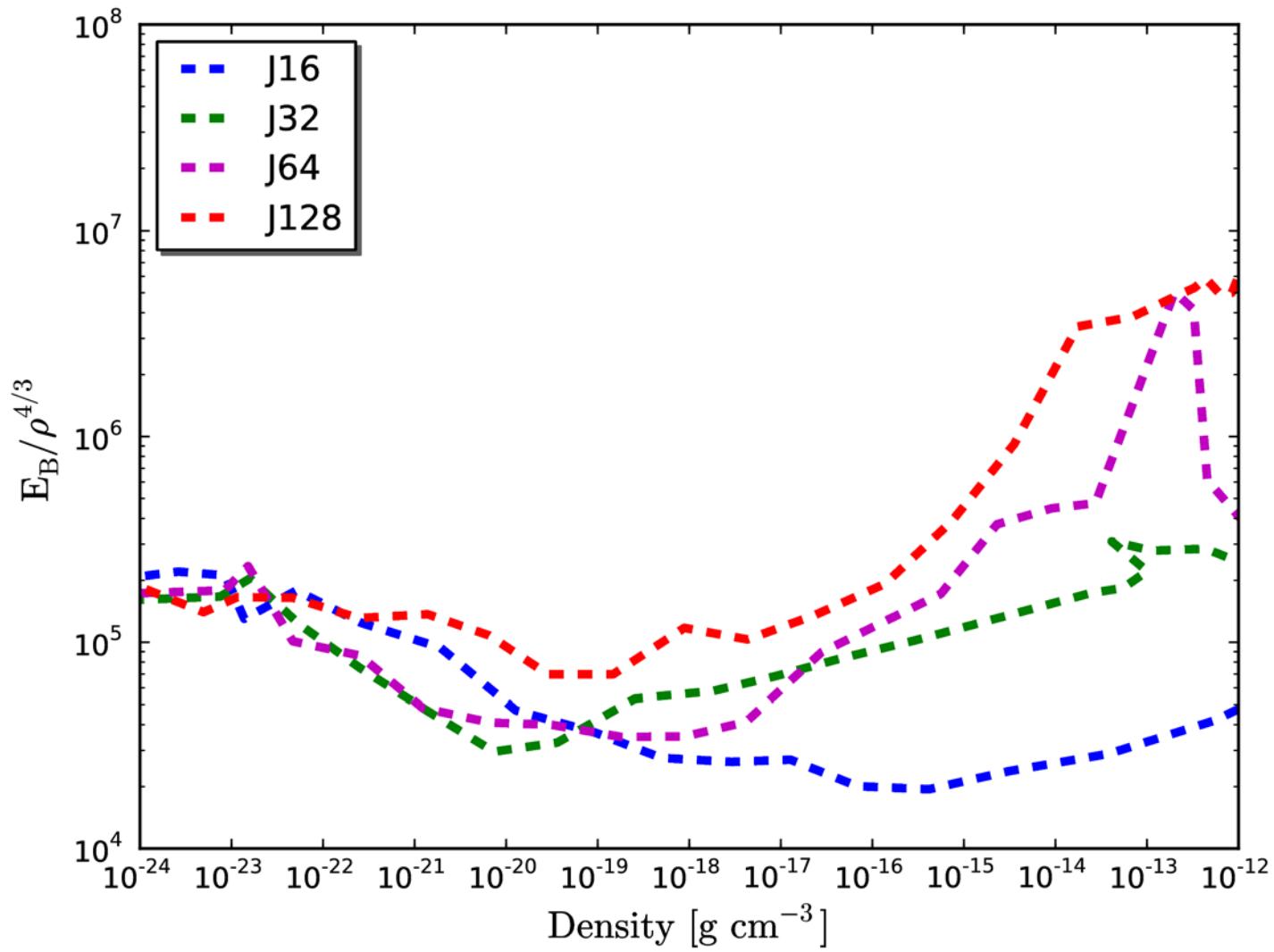
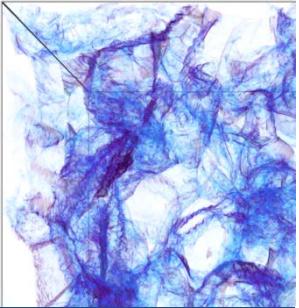


## Application II: The First Stars in the Universe

- **Direct collapse scenario:**
  - dark matter halos of mass  $\sim 10^7 M_{\text{sun}}$
  - fragmentation due to atomic gas cooling
  - collapse produces prestellar cores ( $1000 M_{\text{sun}}$ )
- Might lead to the formation of seed black holes that can grow to **supermassive BHs**
- **Deep zoom-in simulations** with Enzo ([Latif, Schleicher, WS, and Niemeyer 2013](#)):
  - 27 levels of refinement
  - follows collapse down to 0.25 AU
  - MHD runs
  - HD runs: comparison LES to ILES

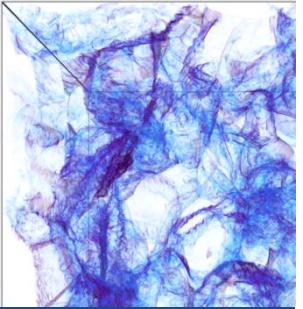






## Magnetic field amplification:

- no amplification at low resolution
- Amplification at high densities for  $\geq 64$  cells per  $\lambda_J$

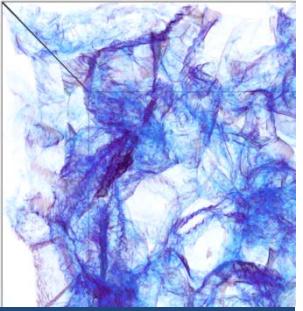


# What's going on here?

- Turbulence is **driven by self-gravity** of the gas
- Energy injection on length scales  $l \geq \lambda_J$
- We barely touch the turbulent cascade in these simulations – would need  $\Delta \ll \lambda_J$
- **Growth rate** of  $B$  due to dynamo can be estimated by (**Schober et al. 2012**):

$$1/\tau_B \sim \frac{V}{\lambda_J} Re^{1/2} \sim \frac{V}{\lambda_J} \left( \frac{\lambda_J}{l_K} \right)^{2/3}$$

- But in ILES that do not resolve the physical dissipation scale, the dynamo is driven from the **smallest resolved length scales**  $l \sim \Delta \gg l_K$



# Subgrid Scale Model for Hydrodynamical Turbulence

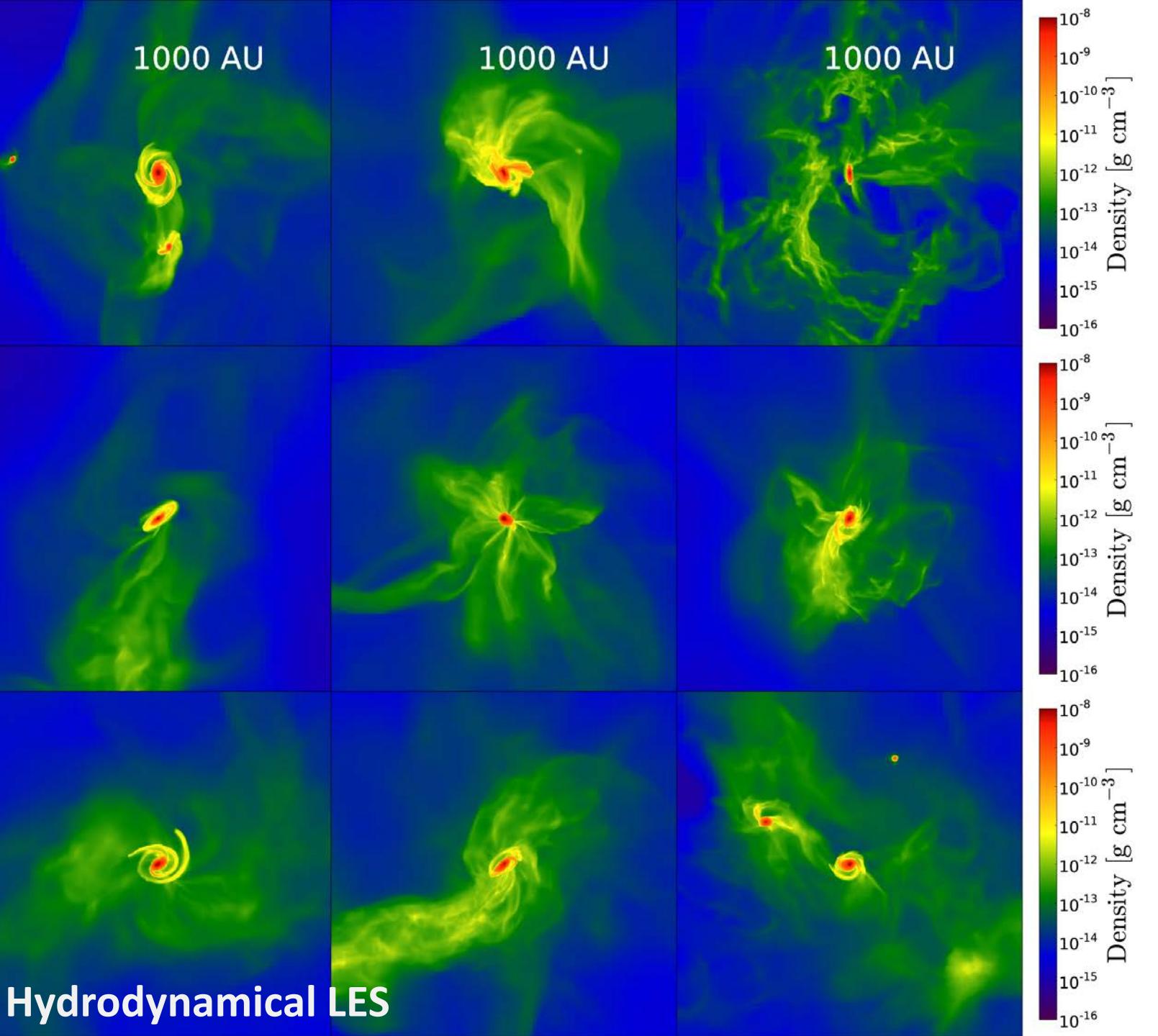
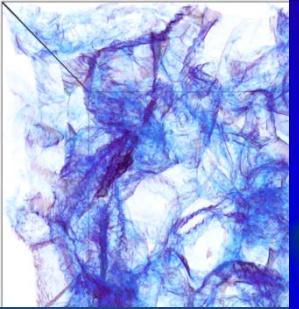
- Based on **Germano (1992)** decomposition
- A priori tests of closure for compressible turbulence (**Schmidt et al. 2006, 2011**)

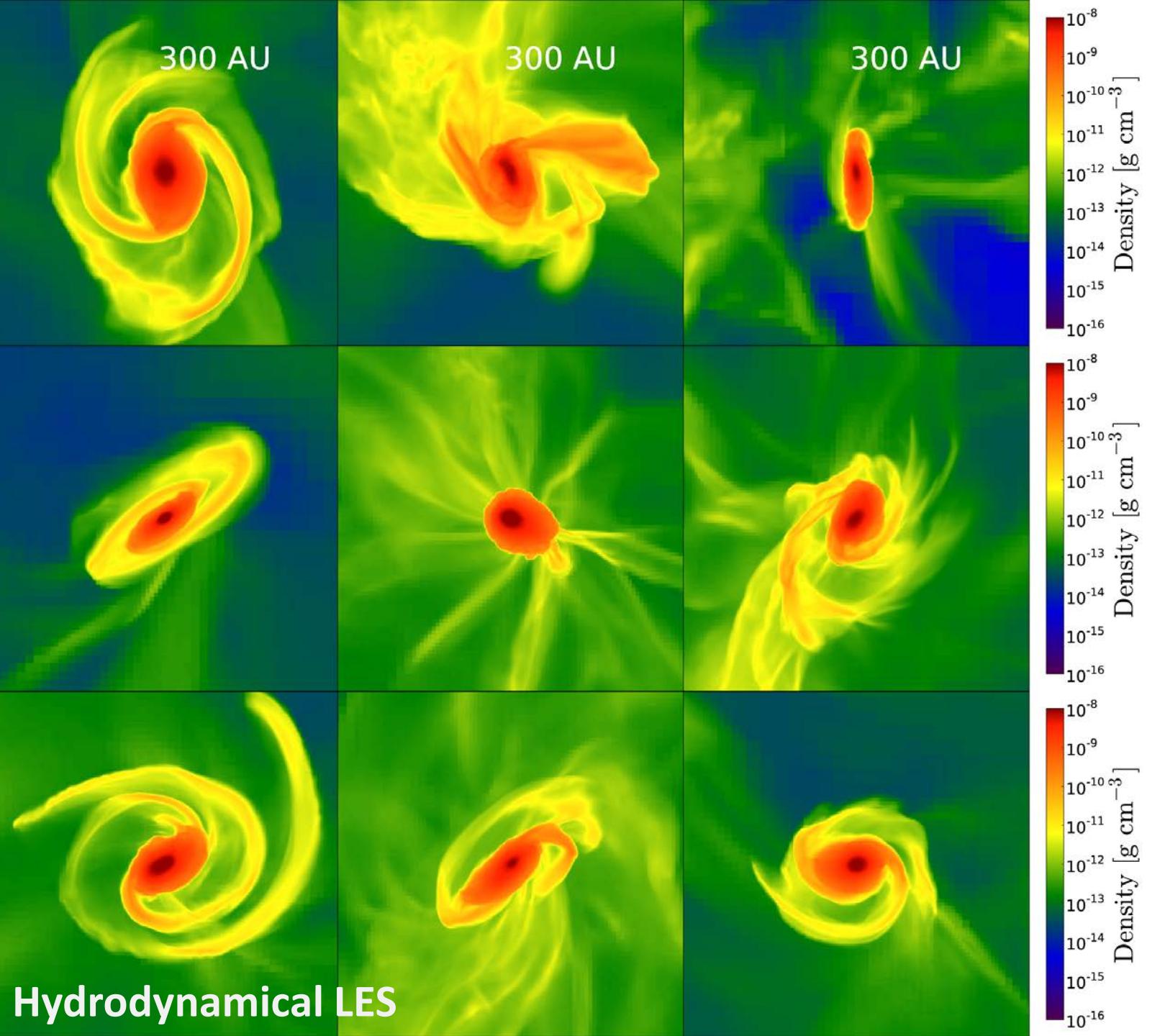
$$\frac{\partial}{\partial t} \rho + \nabla \cdot (\mathbf{u} \rho) = 0$$

$$\frac{\partial}{\partial t} (\rho \mathbf{u}) + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) = - \nabla \left( \underbrace{P + \frac{2}{3} \rho K}_{\text{eff. pressure}} \right) + \underbrace{\nabla \cdot \boldsymbol{\tau}_{\text{sgs}}^*}_{\text{nondiag. stresses}} + \rho (\mathbf{g} + \mathbf{f}_{\text{ext}})$$

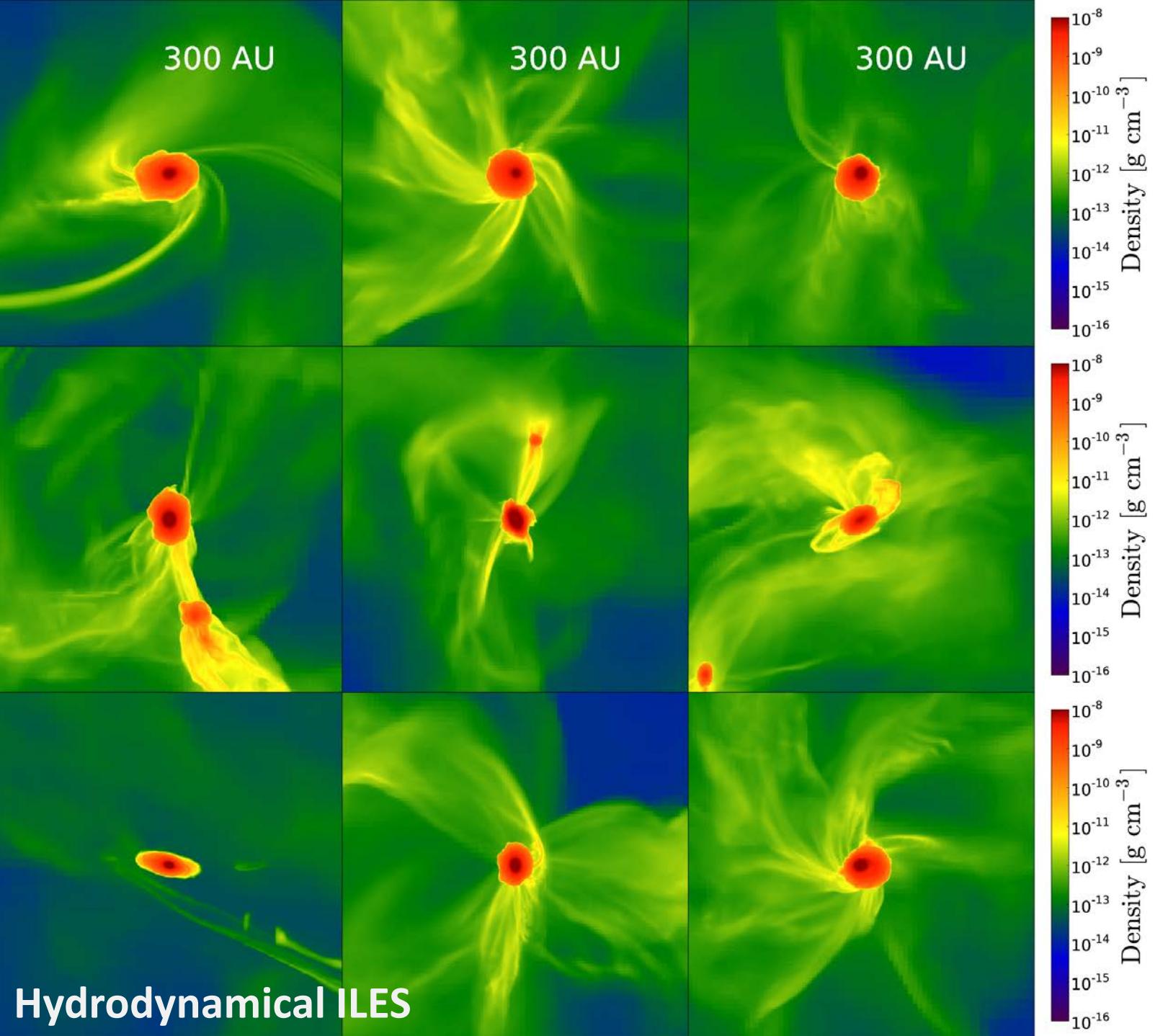
$$\begin{aligned} \frac{\partial}{\partial t} \rho E + \nabla \cdot (\rho \mathbf{u} E) &= - \nabla \cdot \left[ \mathbf{u} \left( P + \frac{2}{3} \rho K \right) \right] + \nabla \cdot (\mathbf{u} \cdot \boldsymbol{\tau}_{\text{sgs}}^*) \\ &\quad + \rho \mathbf{u} \cdot (\mathbf{g} + \mathbf{f}_{\text{ext}}) \underbrace{-\Lambda + \Gamma}_{\text{radiative}} \underbrace{-\Sigma + \rho \epsilon}_{\text{turbulent}} \end{aligned}$$

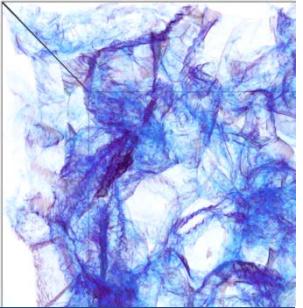
$$\frac{\partial}{\partial t} \rho K + \nabla \cdot (\rho \mathbf{u} K) = \mathfrak{D} + \Sigma - \rho \epsilon$$





## Hydrodynamical ILES





# ILES vs LES

	ILES	LES
diffusivity (SGS)	$\nu_{\text{num}} = \nu_{\text{turb}}?$	e.g. $\nu_{\text{sgs}} = C_v \Delta K^{1/2}$
dissipation	instantaneous (kin. energy to heat)	intermediate reservoir $\rho K$
turbulent pressure (SGS)	none	$P_{\text{sgs}} = \frac{2}{3} \rho K$
dynamo	increases with $1/\Delta$	closure for $\alpha B$
AMR	numerical cooling/heating	energy bookkeeping ( $\frac{1}{2} \rho u^2 \leftrightarrow \rho K$ )